

might be due to packing considerations and one might assume that one chlorine atom is bent  $32^\circ$  from the plane because of an interaction with the pyridine ring. A consideration of contact distances does not show the chlorine atom to be significantly closer when it is placed in the copper-oxygen plane than when it is bent out of the plane. We do not believe that a steric interaction would be sufficiently strong to give the resulting distorted structure. Some careful theoretical

considerations will have to be employed to see if the proposed mechanism is possible.

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## The Crystal Structure of Tris(2-dimethylaminoethyl)aminocobalt(II) Bromide

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The crystal structure of  $\text{Co}(\text{Me}_6\text{tren})\text{Br}_2$  [ $\text{Me}_6\text{tren} = \text{N}(\text{CH}_2\text{CH}_2\text{N}(\text{CH}_3)_2)_3$ ] has been determined by three-dimensional X-ray analysis and refined to a final  $R$  factor of 8.0%. Crystals of the compound are cubic, space group  $\text{P}2_13$ ,  $a = 12.088 \pm 0.004$  Å, with four formula units per unit cell. The structure consists of  $\text{Co}(\text{Me}_6\text{tren})\text{Br}^+$  and  $\text{Br}^-$  ions arranged in a distorted NaCl lattice. The  $\text{Co}(\text{Me}_6\text{tren})\text{Br}^+$  ion has a trigonal bipyramidal structure with symmetry  $\text{C}_3$ . The symmetry of the complex ion is essentially determined by the shape of the ligand molecule. The series of complexes  $\text{M}^{\text{II}}(\text{Me}_6\text{tren})\text{Br}_2$  with  $\text{M}^{\text{II}} = \text{Cr}, \text{Mn}, \text{Fe}, \text{Co}, \text{Ni}, \text{Cu},$  and  $\text{Zn}$  are all isomorphous.

### Introduction

Although several five-coordinated cobalt(II) complexes have been described by X-ray analysis,<sup>1-6</sup> the factors determining the stereochemistry of these complexes are not well understood. From the results it appears that, among the other factors, the geometry of the ligand molecule plays an important role in determining the stereochemistry of the complex.

Ciampolini and Nardi<sup>7</sup> have reported the preparation and the physicochemical properties of a series of high-spin complex compounds with general formula  $\text{M}^{\text{II}}(\text{Me}_6\text{tren})\text{X}_2$ , where  $\text{M}^{\text{II}} = \text{Cr}, \text{Mn}, \text{Fe}, \text{Co}, \text{Ni}, \text{Cu},$  and  $\text{Zn}$  and  $\text{Me}_6\text{tren} = \text{tris}(2\text{-dimethylaminoethyl})\text{amine}, \text{N}(\text{CH}_2\text{CH}_2\text{N}(\text{CH}_3)_2)_3$ , and  $\text{X} = \text{Cl}, \text{Br}, \text{I}, \text{NO}_3,$  and  $\text{ClO}_4$ .

On the basis of magnetic, spectral, and conductivity measurements, these complexes have been formulated as  $[\text{M}(\text{Me}_6\text{tren})\text{X}]\text{X}$ , and a trigonal bipyramidal structure has been proposed for the  $[\text{M}(\text{Me}_6\text{tren})\text{X}]^+$  ion.

All the complexes with the same anion are isomorphous, as shown by the similarity of their X-ray powder patterns. The bromides crystallize in well-formed tetrahedra with dimensions particularly suitable for X-ray investigation. We report here the three-

dimensional X-ray analysis of the complex  $[\text{Co}(\text{Me}_6\text{tren})\text{Br}]\text{Br}$ .

### Experimental Section

**Data Collection.**—Crystals of the complex suitable for the structural analysis were kindly supplied by Dr. Ciampolini. Cell dimensions were determined from a Weissenberg photograph using Fe  $\text{K}\alpha$  radiation ( $\lambda$  1.9373 Å), with the NaCl rotation pattern superimposed on the film for calibration ( $a = 5.6273$  Å). Six  $2\theta$  values were measured and treated by least squares. The results are:  $a = 12.088 \pm 0.004$  Å,  $d_0 = 1.688$  g/cm<sup>3</sup>,  $Z = 4$ ,  $\mu = 139.0$  cm<sup>-1</sup>, space group  $\text{P}2_13$ .

The crystal chosen for the collection of the intensities was a truncated tetrahedron with the edge of about 0.25 mm.

Doubly integrated Weissenberg photographs  $hk0$  through  $hk6$  were taken on a Nonius camera, with the multiple-film, equi-inclination technique. The intensities of the reflections were measured on a Nonius microdensitometer, the intensity being assumed proportional to the density of the spots. The various levels were scaled by means of common reflections. We examined 390 independent reflections, about 95% of the total number in the Fe sphere of reflection; we found 53 too weak to be measured.

Intensities were corrected for Lorentz and polarization effects and for spot elongation on upper levels according to Phillips.<sup>8</sup>

Later, during the determination of the structure, inspection of the observed and calculated structure factors at  $R = 12.1\%$  showed the existence of strong secondary extinction effects. This was confirmed by a plot of  $\ln(I_0/I)$  against  $I_0$  for the strongest reflections, but an attempt to correct all of the reflections on the basis of the plot did not give satisfactory results. It was decided then to continue the refinement, omitting those reflections which appeared to be affected most heavily by extinction. These reflections are 210, 220, 121, 231, 331, 431, 241, 351, 171, 222, and 262.

The atomic scattering factors were taken from the Dirac-

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TABLE I  
 POSITIONAL PARAMETERS, TEMPERATURE FACTORS, AND THEIR ESTIMATED STANDARD DEVIATIONS<sup>a</sup>

Atom	Positional parameter $\times 10^4$			Thermal parameter $\times 10^4$					
	$x/a$	$y/b$	$z/c$	$\beta_{11}$	$\beta_{22}$	$\beta_{33}$	$\beta_{12}$	$\beta_{13}$	$\beta_{23}$
Co	2280 (3)	2280 (3)	2280 (3)	44 (2)	44 (2)	44 (2)	-1 (2)	-1 (2)	-1 (2)
Br(1)	1119 (2)	1119 (2)	1119 (2)	52 (2)	52 (2)	52 (2)	-8 (2)	-8 (2)	-8 (2)
Br(2)	8151 (2)	8151 (2)	8151 (2)	61 (2)	61 (2)	61 (2)	5 (2)	5 (2)	5 (2)
N(1)	3308 (17)	3308 (17)	3308 (17)	60 (16)	60 (16)	60 (16)	10 (15)	10 (15)	10 (15)
N(2)	2752 (16)	1106 (16)	3446 (15)	32 (15)	54 (15)	41 (15)	6 (13)	-4 (13)	-25 (13)
C(1)	3619 (18)	2698 (19)	4292 (18)	49 (16)	64 (17)	8 (17)	2 (15)	-7 (14)	30 (15)
C(2)	3820 (18)	1470 (17)	3993 (18)	21 (15)	46 (16)	47 (16)	23 (14)	-18 (15)	-3 (13)
C(3)	2957 (22)	-54 (23)	2984 (22)	55 (23)	53 (22)	92 (21)	39 (19)	21 (20)	-16 (18)
C(4)	1868 (27)	945 (26)	4294 (25)	103 (25)	62 (25)	77 (24)	-24 (23)	-13 (22)	23 (21)

<sup>a</sup> Anisotropic thermal factors are of the form:  $\exp(-\sum_{i=1}^3 \sum_{j=1}^3 h_i h_j \beta_{ij})$ . Estimated standard deviations are given in parentheses.

Slater calculations of Cromer and Waber.<sup>9</sup> No correction was made for anomalous dispersion or for absorption effects.

**Structure Determination.**—The positions of the heavy atoms were determined by means of a three-dimensional Patterson synthesis. There are eight bromine and four cobalt atoms in the unit cell. Examination of the equivalent positions of the space group P2<sub>1</sub>3 shows that these atoms must lie on the threefold axes, *i.e.*, in 4a:  $x, x, x; 1/2 + x, 1/2 - x, \bar{x}; \bar{x}, 1/2 + x, 1/2 - x; 1/2 - x, \bar{x}, 1/2 + x$ . Each set of equivalent atoms generates a set of 12 peaks in the vector space lying in the three Harker sections having, respectively,  $u = 1/2, v = 1/2, w = 1/2$ . On the other hand, two different sets of equivalent atoms in 4a give rise to 32 vectors, eight of which lie along the threefold axes. One of the bromine atoms was located from one of the two Harker sets which showed the highest peaks. The second bromine atom and the cobalt atom were successively located from the peaks lying along the threefold axes. These positions were then checked for all of the highest peaks in the Patterson synthesis.

The Fourier synthesis phased on the positions of the heavy atoms showed clearly all the other nonhydrogen atoms. At this point the *R* factor was 13.6%, but some of the bond lengths between light atoms were rather improbable. *R* is throughout defined as  $\Sigma ||F_o| - |F_c|| / \Sigma |F_o|$ , where the sums are over the independent, observed reflections.

A  $\Delta F$  synthesis calculated at this point did not give any decisive improvement. Refinement was then continued by means of block-diagonal, least-squares analysis, with individual isotropic temperature factors for the atoms. Reflections were weighted according to the function:<sup>10</sup>  $w = 1/(a + F_o + cF_o^2)$ , with  $a = 30.0$  and  $c = 0.02$  Å. Unobserved reflections were not introduced in this series of calculations. A set of three cycles reduced the *R* factor to 12.1% but the C-C bond length still appeared to be rather long (1.69 Å).

It was then decided to continue the refinement with the full-matrix, least-squares program of Busing and Levy, adapted for the 7090 by Stewart.<sup>11</sup>

Unobserved reflections were introduced in this set of calculations with the minimum observable intensity. Reflections affected by extinction were omitted as mentioned above. The weighting scheme used was as follows:  $\sqrt{w} = 0$  for unobserved reflections with  $F_o \leq F_{min}$ ;  $\sqrt{w} = 1$  for unobserved reflections with  $F_o > F_{min}$  and for observed reflections with  $F_o \leq 60$ ;  $\sqrt{w} = 60/F_o$  for observed reflections with  $F_o > 60$ .

Three cycles with individual isotropic temperature factors reduced the *R* factor to 9.1%. The C-C bond length showed a decisive improvement (1.59 Å).

At this point the hydrogen atoms were introduced in calculated positions with a thermal parameter  $B = 4.0$  Å<sup>2</sup>, and their coordinates were not refined. Individual anisotropic temperature factors were applied to all the nonhydrogen atoms. One more

cycle of least squares gave a final *R* factor of 8.0% and for the C-C bond a satisfactory value of 1.55 Å. It should be noted that in this cycle the temperature factor of C(1) went nonpositive definite, probably owing to the uncorrected absorption.

Table I reports the final atomic parameters and the temperature factors with their estimated standard deviations. Standard deviations of the parameters were calculated from the expression

$$\sigma_i = \sqrt{\frac{(\Sigma w \Delta^2) a_{ii}}{n - s}}$$

where  $a_{ii}$  is the diagonal term of the inverse least-squares matrix,  $n$  is the number of observations, and  $s$  the number of parameters. Table II reports the coordinates of the hydrogen atoms and Table III the observed and calculated structure factors.

 TABLE II  
 COORDINATES OF HYDROGEN ATOMS

Atom	$10^4x/a$	$10^4y/b$	$10^4z/c$
H'(1)	4343	3066	4646
H''(1)	2951	2760	4889
H'(2)	4515	1382	3431
H''(2)	3989	0973	4715
H'(3)	2197	-0290	2584
H''(3)	3612	-0030	2380
H'''(3)	3160	-0609	3625
H'(4)	1088	0685	3926
H''(4)	2104	0293	4875
H'''(4)	1716	1682	4761

All of the calculations, with the exception of the full-matrix, least-squares cycles, were performed with a set of programs written for the IBM 1620 by Albano, Bellon, Pompa, and Scatturin.<sup>12</sup>

## Discussion

The structure of Co(Me<sub>6</sub>tren)Br<sub>2</sub> consists of Co-(Me<sub>6</sub>tren)Br<sup>+</sup> and Br<sup>-</sup> ions arranged in a distorted NaCl lattice. Each cation is surrounded by three Br<sup>-</sup> ions at 5.65 Å and three at 6.67 Å, forming a distorted octahedron.

The Co(Me<sub>6</sub>tren)Br<sup>+</sup> ion (Figure 1) has a trigonal bipyramidal structure with symmetry C<sub>3</sub>. The cobalt atom is five-coordinated by the four nitrogen atoms of the amine and by one bromine atom. Bond lengths and angles in the ion with their estimated standard deviations are reported in Table IV.

The average Co-N distance of 2.11 Å is in good agree-

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TABLE III  
OBSERVED AND CALCULATED STRUCTURE FACTORS<sup>a</sup>

H, I, 0	Observed	Calculated	H, I, 0	Observed	Calculated
1 000 000 0	1.0000	1.0000	1 000 000 0	1.0000	1.0000
2 000 000 0	0.9999	0.9999	2 000 000 0	0.9999	0.9999
3 000 000 0	0.9997	0.9997	3 000 000 0	0.9997	0.9997
4 000 000 0	0.9993	0.9993	4 000 000 0	0.9993	0.9993
5 000 000 0	0.9987	0.9987	5 000 000 0	0.9987	0.9987
6 000 000 0	0.9979	0.9979	6 000 000 0	0.9979	0.9979
7 000 000 0	0.9969	0.9969	7 000 000 0	0.9969	0.9969
8 000 000 0	0.9957	0.9957	8 000 000 0	0.9957	0.9957
9 000 000 0	0.9943	0.9943	9 000 000 0	0.9943	0.9943
10 000 000 0	0.9927	0.9927	10 000 000 0	0.9927	0.9927
11 000 000 0	0.9909	0.9909	11 000 000 0	0.9909	0.9909
12 000 000 0	0.9889	0.9889	12 000 000 0	0.9889	0.9889
13 000 000 0	0.9867	0.9867	13 000 000 0	0.9867	0.9867
14 000 000 0	0.9843	0.9843	14 000 000 0	0.9843	0.9843
15 000 000 0	0.9817	0.9817	15 000 000 0	0.9817	0.9817
16 000 000 0	0.9789	0.9789	16 000 000 0	0.9789	0.9789
17 000 000 0	0.9759	0.9759	17 000 000 0	0.9759	0.9759
18 000 000 0	0.9727	0.9727	18 000 000 0	0.9727	0.9727
19 000 000 0	0.9693	0.9693	19 000 000 0	0.9693	0.9693
20 000 000 0	0.9657	0.9657	20 000 000 0	0.9657	0.9657
21 000 000 0	0.9619	0.9619	21 000 000 0	0.9619	0.9619
22 000 000 0	0.9579	0.9579	22 000 000 0	0.9579	0.9579
23 000 000 0	0.9537	0.9537	23 000 000 0	0.9537	0.9537
24 000 000 0	0.9493	0.9493	24 000 000 0	0.9493	0.9493
25 000 000 0	0.9447	0.9447	25 000 000 0	0.9447	0.9447
26 000 000 0	0.9399	0.9399	26 000 000 0	0.9399	0.9399
27 000 000 0	0.9349	0.9349	27 000 000 0	0.9349	0.9349
28 000 000 0	0.9297	0.9297	28 000 000 0	0.9297	0.9297
29 000 000 0	0.9243	0.9243	29 000 000 0	0.9243	0.9243
30 000 000 0	0.9187	0.9187	30 000 000 0	0.9187	0.9187
31 000 000 0	0.9129	0.9129	31 000 000 0	0.9129	0.9129
32 000 000 0	0.9069	0.9069	32 000 000 0	0.9069	0.9069
33 000 000 0	0.9007	0.9007	33 000 000 0	0.9007	0.9007
34 000 000 0	0.8943	0.8943	34 000 000 0	0.8943	0.8943
35 000 000 0	0.8877	0.8877	35 000 000 0	0.8877	0.8877
36 000 000 0	0.8809	0.8809	36 000 000 0	0.8809	0.8809
37 000 000 0	0.8739	0.8739	37 000 000 0	0.8739	0.8739
38 000 000 0	0.8667	0.8667	38 000 000 0	0.8667	0.8667
39 000 000 0	0.8593	0.8593	39 000 000 0	0.8593	0.8593
40 000 000 0	0.8517	0.8517	40 000 000 0	0.8517	0.8517
41 000 000 0	0.8439	0.8439	41 000 000 0	0.8439	0.8439
42 000 000 0	0.8359	0.8359	42 000 000 0	0.8359	0.8359
43 000 000 0	0.8277	0.8277	43 000 000 0	0.8277	0.8277
44 000 000 0	0.8193	0.8193	44 000 000 0	0.8193	0.8193
45 000 000 0	0.8107	0.8107	45 000 000 0	0.8107	0.8107
46 000 000 0	0.8019	0.8019	46 000 000 0	0.8019	0.8019
47 000 000 0	0.7929	0.7929	47 000 000 0	0.7929	0.7929
48 000 000 0	0.7837	0.7837	48 000 000 0	0.7837	0.7837
49 000 000 0	0.7743	0.7743	49 000 000 0	0.7743	0.7743
50 000 000 0	0.7647	0.7647	50 000 000 0	0.7647	0.7647
51 000 000 0	0.7549	0.7549	51 000 000 0	0.7549	0.7549
52 000 000 0	0.7449	0.7449	52 000 000 0	0.7449	0.7449
53 000 000 0	0.7347	0.7347	53 000 000 0	0.7347	0.7347
54 000 000 0	0.7243	0.7243	54 000 000 0	0.7243	0.7243
55 000 000 0	0.7137	0.7137	55 000 000 0	0.7137	0.7137
56 000 000 0	0.7029	0.7029	56 000 000 0	0.7029	0.7029
57 000 000 0	0.6919	0.6919	57 000 000 0	0.6919	0.6919
58 000 000 0	0.6807	0.6807	58 000 000 0	0.6807	0.6807
59 000 000 0	0.6693	0.6693	59 000 000 0	0.6693	0.6693
60 000 000 0	0.6577	0.6577	60 000 000 0	0.6577	0.6577
61 000 000 0	0.6459	0.6459	61 000 000 0	0.6459	0.6459
62 000 000 0	0.6339	0.6339	62 000 000 0	0.6339	0.6339
63 000 000 0	0.6217	0.6217	63 000 000 0	0.6217	0.6217
64 000 000 0	0.6093	0.6093	64 000 000 0	0.6093	0.6093
65 000 000 0	0.5967	0.5967	65 000 000 0	0.5967	0.5967
66 000 000 0	0.5839	0.5839	66 000 000 0	0.5839	0.5839
67 000 000 0	0.5709	0.5709	67 000 000 0	0.5709	0.5709
68 000 000 0	0.5577	0.5577	68 000 000 0	0.5577	0.5577
69 000 000 0	0.5443	0.5443	69 000 000 0	0.5443	0.5443
70 000 000 0	0.5307	0.5307	70 000 000 0	0.5307	0.5307
71 000 000 0	0.5169	0.5169	71 000 000 0	0.5169	0.5169
72 000 000 0	0.5029	0.5029	72 000 000 0	0.5029	0.5029
73 000 000 0	0.4887	0.4887	73 000 000 0	0.4887	0.4887
74 000 000 0	0.4743	0.4743	74 000 000 0	0.4743	0.4743
75 000 000 0	0.4597	0.4597	75 000 000 0	0.4597	0.4597
76 000 000 0	0.4449	0.4449	76 000 000 0	0.4449	0.4449
77 000 000 0	0.4299	0.4299	77 000 000 0	0.4299	0.4299
78 000 000 0	0.4147	0.4147	78 000 000 0	0.4147	0.4147
79 000 000 0	0.3993	0.3993	79 000 000 0	0.3993	0.3993
80 000 000 0	0.3837	0.3837	80 000 000 0	0.3837	0.3837
81 000 000 0	0.3679	0.3679	81 000 000 0	0.3679	0.3679
82 000 000 0	0.3519	0.3519	82 000 000 0	0.3519	0.3519
83 000 000 0	0.3357	0.3357	83 000 000 0	0.3357	0.3357
84 000 000 0	0.3193	0.3193	84 000 000 0	0.3193	0.3193
85 000 000 0	0.3027	0.3027	85 000 000 0	0.3027	0.3027
86 000 000 0	0.2859	0.2859	86 000 000 0	0.2859	0.2859
87 000 000 0	0.2689	0.2689	87 000 000 0	0.2689	0.2689
88 000 000 0	0.2517	0.2517	88 000 000 0	0.2517	0.2517
89 000 000 0	0.2343	0.2343	89 000 000 0	0.2343	0.2343
90 000 000 0	0.2167	0.2167	90 000 000 0	0.2167	0.2167
91 000 000 0	0.1989	0.1989	91 000 000 0	0.1989	0.1989
92 000 000 0	0.1809	0.1809	92 000 000 0	0.1809	0.1809
93 000 000 0	0.1627	0.1627	93 000 000 0	0.1627	0.1627
94 000 000 0	0.1443	0.1443	94 000 000 0	0.1443	0.1443
95 000 000 0	0.1257	0.1257	95 000 000 0	0.1257	0.1257
96 000 000 0	0.1069	0.1069	96 000 000 0	0.1069	0.1069
97 000 000 0	0.0879	0.0879	97 000 000 0	0.0879	0.0879
98 000 000 0	0.0687	0.0687	98 000 000 0	0.0687	0.0687
99 000 000 0	0.0493	0.0493	99 000 000 0	0.0493	0.0493
100 000 000 0	0.0297	0.0297	100 000 000 0	0.0297	0.0297

<sup>a</sup> Columns are *l* index, 10*F*<sub>0</sub>, 10*F*<sub>1</sub>, and  $\alpha$  (millicycles). Unobserved reflections are marked by L; reflections considered affected by extinction are marked by E.

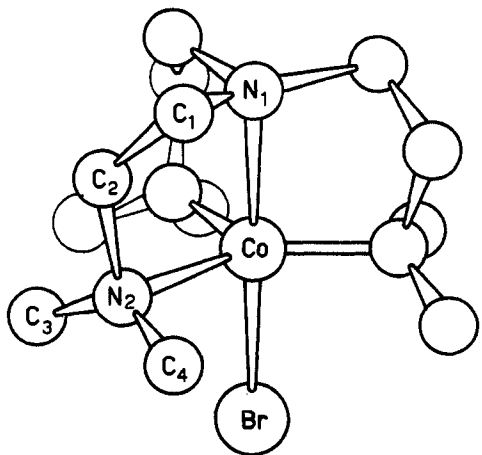


Figure 1.—A perspective drawing of the Co(Me<sub>6</sub>tren)Br<sup>+</sup> ion.

TABLE IV  
DISTANCES AND ANGLES IN THE Co(Me<sub>6</sub>tren)Br<sup>+</sup> ION WITH THEIR ESTIMATED STANDARD DEVIATIONS<sup>a</sup>

Distance, Å	Angle, deg
Co-N(1) 2.151 0.021	N(1)-Co-N(2) 81.1 0.8
Co-N(2) 2.080 0.018	N(2)-Co-Br(1) 98.9 0.5
Co-Br(1) 2.431 0.004	N(2)-Co-N(2')
N(1)-C(1) 1.449 0.029	N(1)-C(1)-C(2) 109.3 1.4
C(1)-C(2) 1.547 0.031	N(1)-C(1)-C(2) 109.7 1.8
C(2)-N(2) 1.515 0.028	C(1)-C(2)-N(2) 104.3 1.7
N(2)-C(3) 1.529 0.033	C(2)-N(2)-C(3) 106.7 1.7
N(2)-C(4) 1.492 0.036	C(2)-N(2)-C(4) 110.4 1.8
Br(1)-C(3) 3.468 0.026	C(3)-N(2)-C(4) 104.3 1.9
Br(1)-C(4) 3.950 0.030	Co-N(2)-C(3) 115.0 1.4
Br(1)-N(2) 3.437 0.018	Co-N(2)-C(4) 111.0 1.6
N(2)-N(2') 3.558 0.026	Co-N(2)-C(2) 109.3 1.3

<sup>a</sup> Primes refer to atoms related to reference atom by the three-fold axis.

ment with the analogous 2.12 Å reported for (1,3-bis-(2'-pyridyl)-2,3-diaza-1-propene)cobalt(II) chloride<sup>4</sup> and with the 2.13 Å of Co(Me<sub>3</sub>dien)Cl<sub>2</sub><sup>3</sup> [Me<sub>3</sub>dien = (CH<sub>3</sub>)<sub>2</sub>NCH<sub>2</sub>CH<sub>2</sub>N(CH<sub>3</sub>)CH<sub>2</sub>CH<sub>2</sub>N(CH<sub>3</sub>)<sub>2</sub>]. It should be noted that the axial Co-N distance is about three standard deviations longer than the equatorial Co-N distances, but the difference is hardly significant since standard errors in Table IV are probably underestimated. The comparison with the structures of the isomorphous complexes of the series, in progress in our laboratory, will further elucidate this question.

The Co-Br distance of 2.43 Å is in good agreement with the sum (2.43 Å) of the Pauling covalent radii of the cobalt (octahedral) and of the bromine. Since the bond cannot be purely covalent, owing to the different electronegativity of the two elements, this seems to indicate the existence of  $\pi$  bonding between the metal and the halogen atom, which would contribute to the stability of the complex.

The N(1)-Co-N(2) angle of 81.1° appears essentially imposed by the stereochemistry of the ligand molecule, although this value is somewhat smaller than those found in ethylenediamine metal chelates.<sup>13</sup> This causes the metal atom to lie 0.32 Å below the equatorial plane of the three nitrogen atoms. The angles formed by the equatorial nitrogen atoms at the cobalt are 117.6°. Bond lengths and angles in the ligand molecule appear to be normal (Table IV).

Gillespie<sup>14</sup> has discussed the main factors which determine the configuration of five-coordinate complexes. Although the electrostatic repulsions between the ligands favor the trigonal bipyramidal configuration, crystal field stabilization energy favors the square-pyramidal configuration, and this effect should predominate in complexes with essentially ionic binding. Anyhow, in most cases, both effects will be comparable and the resultant structure will be intermediate between the two configurations, or other factors, such as the shape of the ligand molecule or packing effects in the solid state, may determine the structure of the complex.

Of the reported structures of five-coordinated cobalt(II) complexes, three<sup>1,2,4</sup> have been described as more or less distorted square pyramids, and three<sup>3,5,6</sup> are intermediate between the two configurations. Co(Me<sub>6</sub>tren)Br<sub>2</sub> is then the first five-coordinate cobalt(II) complex having a trigonal bipyramidal structure with C<sub>3</sub> symmetry. Undoubtedly the symmetry of the complex cation is determined by the particular shape of the ligand molecule.

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